Hadron-Nucleus Interactions

Beginners' FLUKA Course

The FLUKA hadronic Models

Hadron-nucleus: PEANUT

Elastic,exchange
Phase shifts
data, eikonal

P<3-5GeV/c
Resonance prod
and decay

low E
π,K
Special

High Energy
DPM
hadronization

Sophisticated
G-Intranuclear Cascade

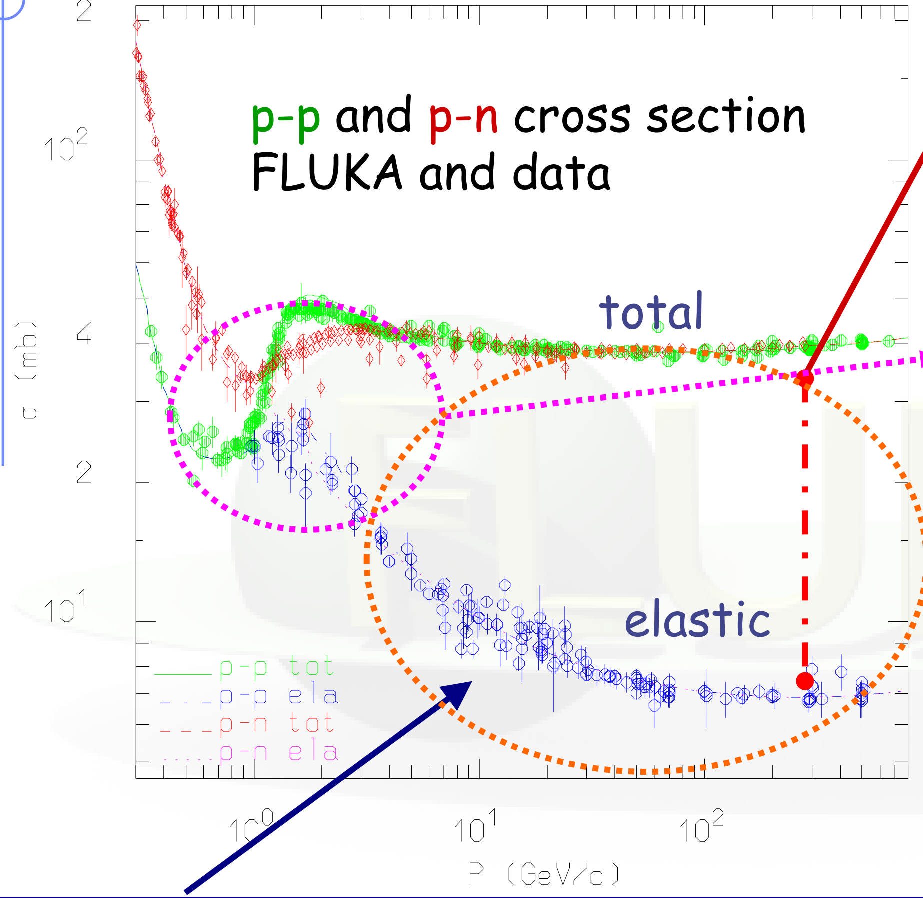
Gradual onset of
Glauber-Gribov multiple
interactions

Preequilibrium

Coalescence

Evaporation/Fission/Fermi break-up
γ deexcitation

Hadron-nucleon interaction models



Particle production interactions: two kinds of models

Those based on "resonance" production and decays, cover the energy range up to 3–5 GeV

Those based on quark/parton string models, which provide reliable results up to several tens of TeV

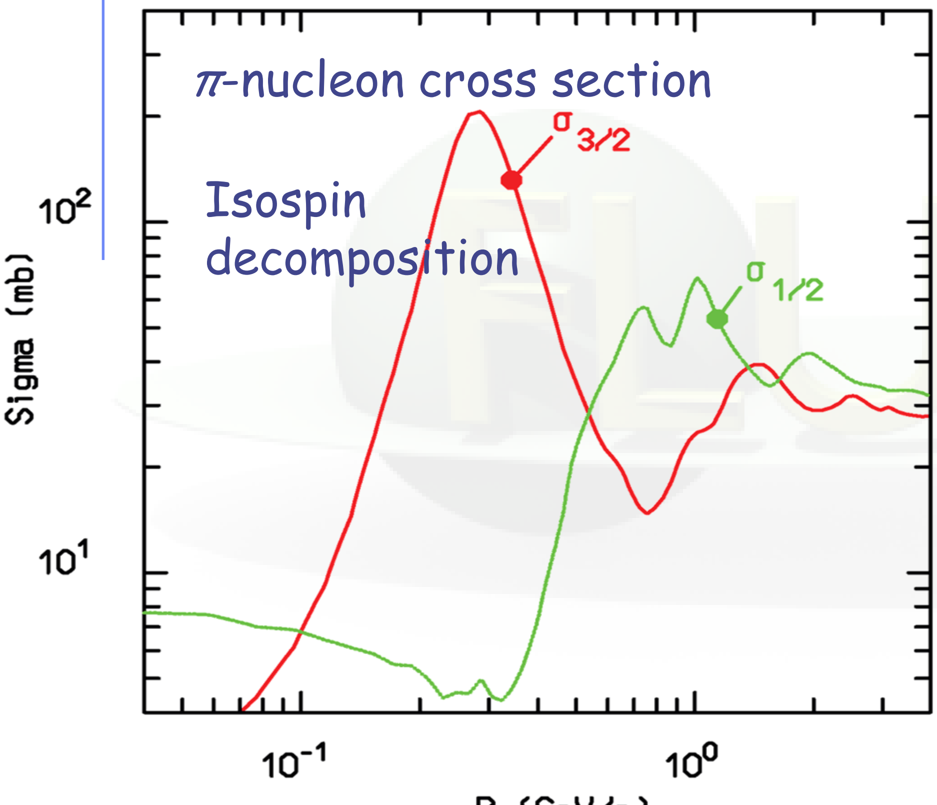
Elastic, charge exchange and strangeness exchange reactions:

- Available phase-shift analysis and/or fits of experimental differential data
- At high energies, standard eikonal approximations are used

Nonelastic hN interactions at intermediate energies

- $N_1 + N_2 \rightarrow N'_1 + N'_2 + \pi$ threshold at 290 MeV, important above 700 MeV,
- $\pi + N \rightarrow \pi' + \pi'' + N'$ opens at 170 MeV.

Anti-nucleon -nucleon open at rest !



Dominance of the Δ resonance and of the $N*$ resonances

$\rightarrow$ isobar model

$\rightarrow$ all reactions proceed through an intermediate state containing at least one resonance.

Resonance energies, widths, cross sections, branching ratios from data and conservation laws, whenever possible. Inferred from inclusive cross sections when needed

$$N_1 + N_2 \rightarrow N''_1 + \Delta(1232) \rightarrow N'_1 + N'_2 + \pi$$

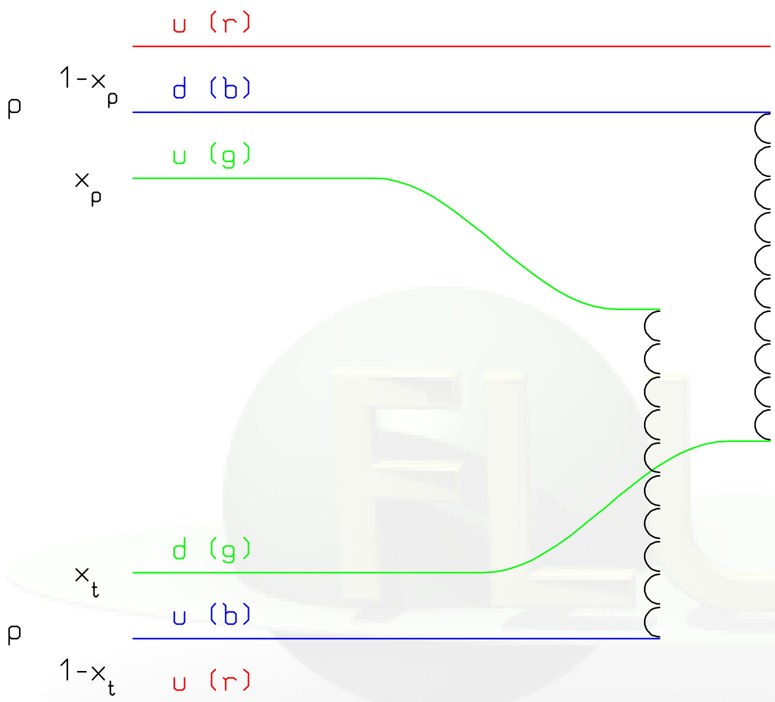
$$\pi + N \rightarrow \Delta(1600) \rightarrow \pi' + \Delta(1232) \rightarrow \pi' + \pi'' + N'$$

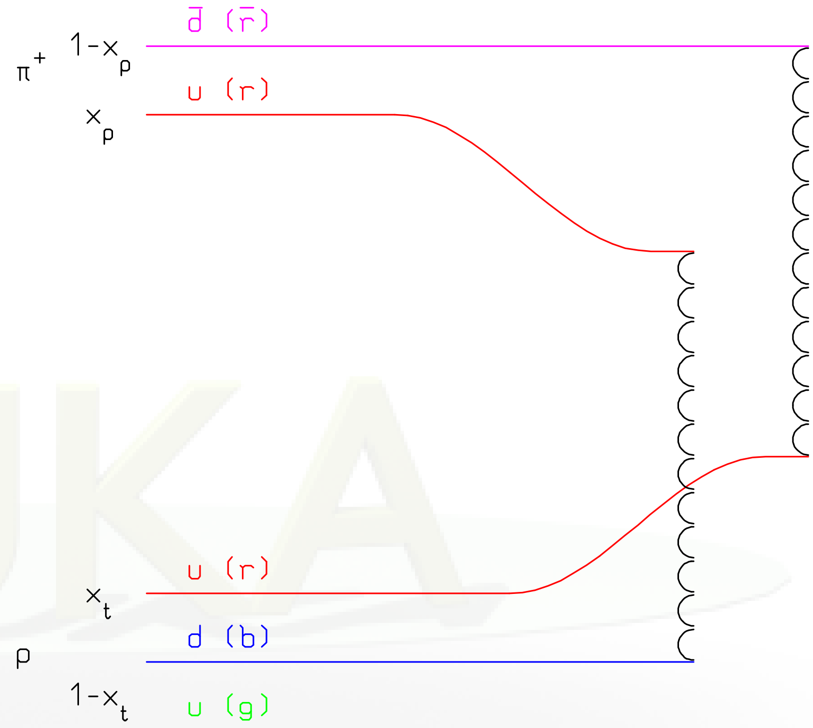
$$N_1 + N_2 \rightarrow \Delta_1(1232) + \Delta_2(1232) \rightarrow N'_1 + \pi_1 + N'_2 + \pi_2$$

Inelastic hN at high energies: (DPM, QGSM, ...)

- Problem: "soft" interactions → QCD perturbation theory cannot be applied.
- Interacting strings (quarks held together by the gluon-gluon interaction into the form of a string)
- Interactions treated in the Reggeon-Pomeron framework
- each of the two hadrons splits into 2 colored partons → combination into 2 colourless chains → 2 back-to-back jets
- each jet is then hadronized into physical hadrons

Hadron-hadron collisions: chain examples





Leading two-chain diagram in DPM for p-p scattering. The color (red, blue, and green) and quark combination shown in the figure is just one of the allowed possibilities

Leading two-chain diagram in DPM for π^+-p scattering. The color (red, blue, and green) and quark combination shown in the figure is just one of the allowed possibilities

DPM and hadronization

from DPM:

- Number of chains
- Chain composition
- Chain energies and momenta
- Diffractive events

Almost No Freedom

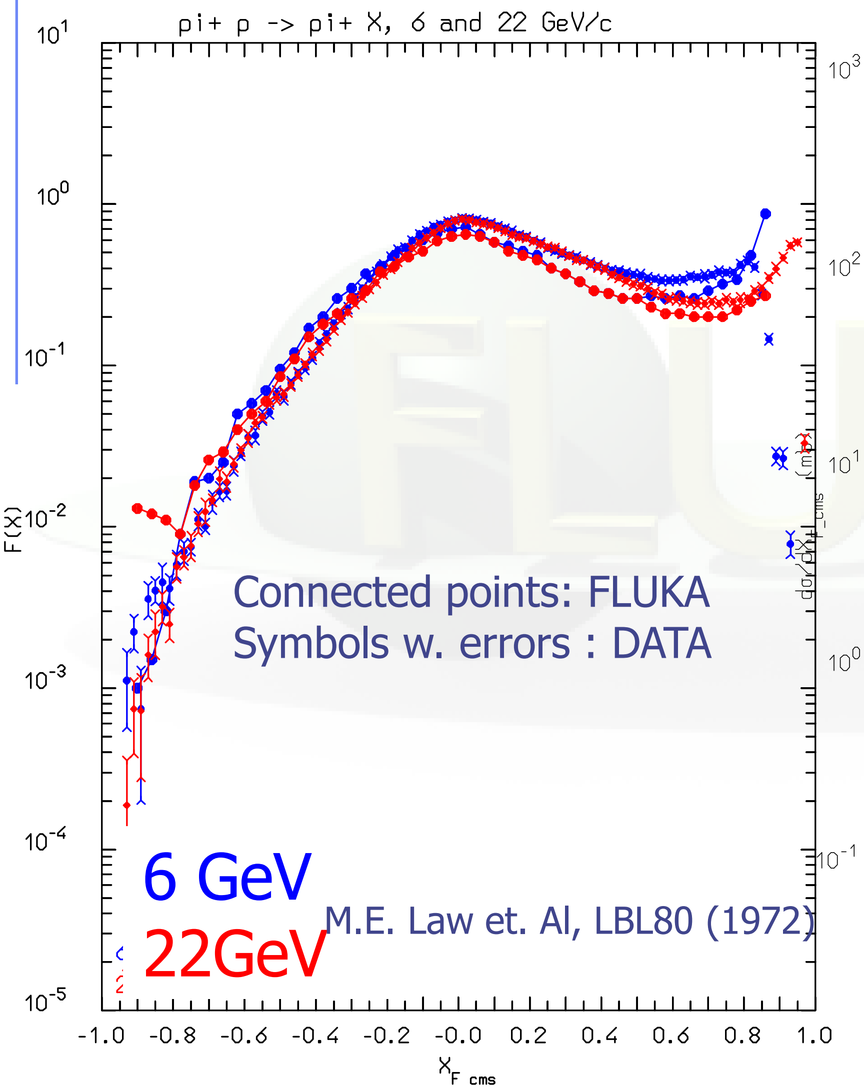
Chain hadronization

- Assumes chain universality
- Fragmentation functions from hard processes and e^+e^- scattering
- Transverse momentum from uncertainty considerations
- Mass effects at low energies

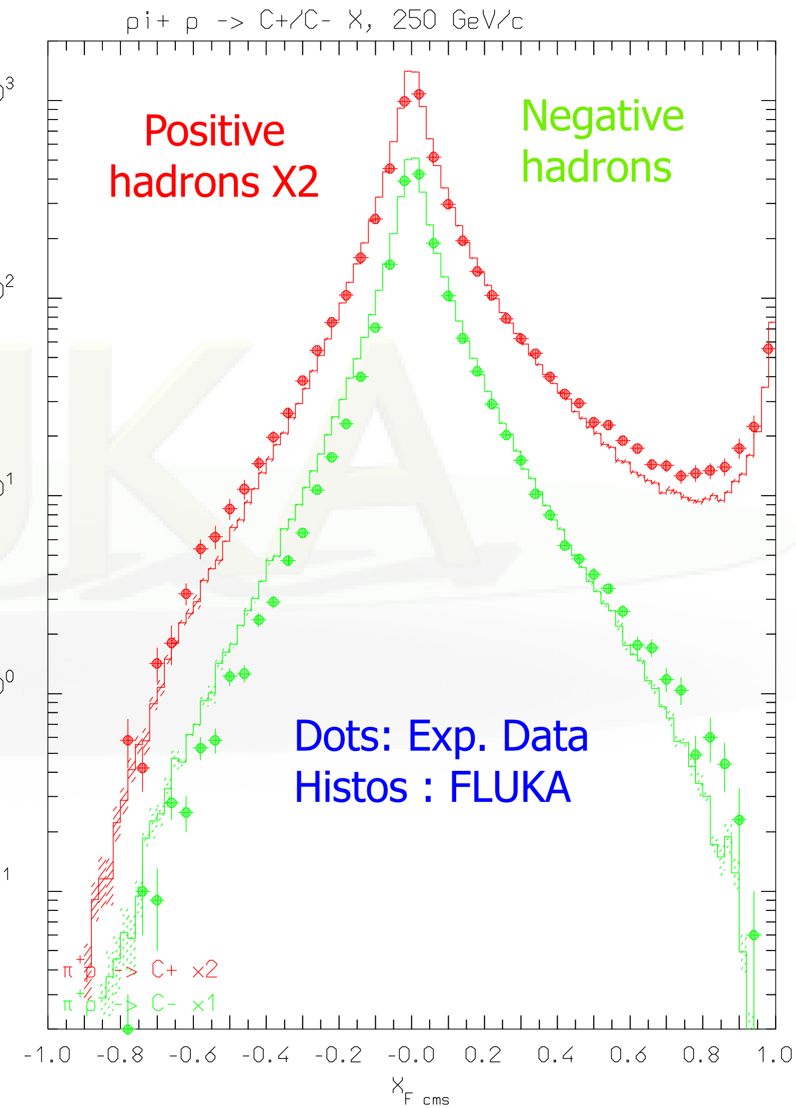
The same functions and (few) parameters for all reactions and energies

Inelastic hN interactions: examples

$\pi^+ + p \rightarrow \pi^+ + X$ (6 & 22 GeV/c)



$\pi^+ + p \rightarrow Ch^+/Ch^- + X$ (250 GeV/c)



Nuclear interactions in PEANUT:

Target nucleus description (density, Fermi motion, etc)

↓

Glauber-Gribov cascade with formation zone

↓

Generalized IntraNuclear cascade

↓

Preequilibrium stage with current exciton configuration and excitation energy
(all non-nucleons emitted/decayed + all nucleons below 30-100 MeV)

↓

Evaporation/Fragmentation/Fission model

↓

γ deexcitation

t (s)

10^{-23}

10^{-22}

10^{-20}

10^{-16}

Nucleon Fermi Motion

- Fermi gas model: Nucleons = Non-interacting Constrained Fermions

  Momentum distribution $\propto \frac{dN}{dk} = \frac{|k|^2}{2\pi^2}$

  for *k* up to a (local) Fermi momentum $k_F(r)$ given by

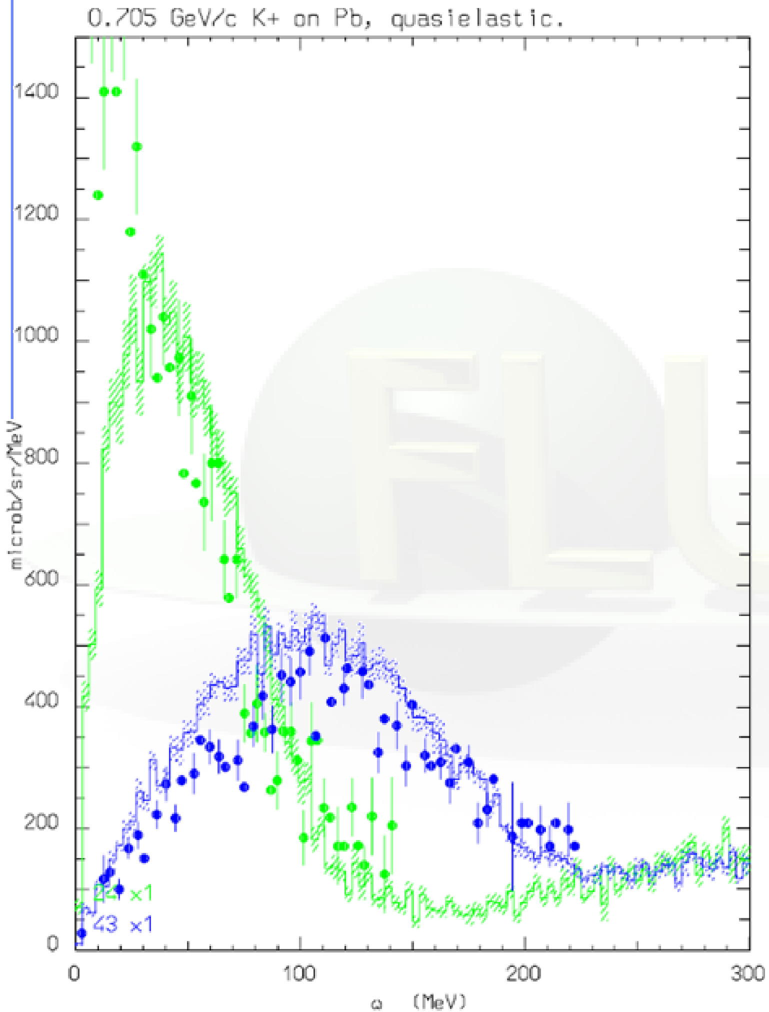
$$k_F(r) = \left[3\pi^2 \rho_N(r)\right]^{\frac{1}{3}}$$

  The Fermi energy ($k_F \approx$ 1.36 fm, $P_F \approx$ 260 MeV/c, $E_F \approx$ 35 MeV, at nuclear max. density) is customarily used in building a self-consistent Nuclear Potential

$\Downarrow$

Depth of the potential well $\equiv$ Fermi Energy + Nuclear Binding Energy

Positive kaons as a probe of Fermi motion



$K^+ \quad K^0$

No low mass S=1 baryons $\rightarrow$
weak K^+N interaction
only elastic and ch. exch. up tc
≈ 800 MeV/c

$(K^+, K^{+\prime})$ on Pb vs residual excita-
tion, 705 MeV/c, at $24°$ and $43°$.
Histo: FLUKA, dots: data (Phys
Rev. C51, 669 (1995))

On free nucleon: recoil energy :
43 MeV at $24°$, 117 MeV at $43°$.

(Generalized) IntraNuclear Cascade

- Primary and secondary particles moving in the nuclear medium
- Target nucleons motion and nuclear well according to the Fermi gas model
- Interaction probability

  σ_{free} + Fermi motion × $\rho(r)$ + exceptions (ex. π)
- Glauber cascade at higher energies
- Classical trajectories (+) nuclear mean potential (resonant for π)
- Curvature from nuclear potential $\rightarrow$ refraction and reflection
- Interactions are incoherent and uncorrelated
- Interactions in projectile-target nucleon CMS $\rightarrow$ Lorentz boosts
- Multibody absorption for π, μ^-, K^-
- Quantum effects (Pauli, formation zone, correlations…)
- Exact conservation of energy, momenta and all additive quantum numbers, including nuclear recoil

hA at high energies: Glauber-Gribov cascade with formation zone

- Glauber cascade
  - Quantum mechanical method to compute Elastic, Quasi-elastic and Absorption hA cross sections from Free hadron-nucleon scattering + nuclear ground state
  - Multiple Collision expansion of the scattering amplitude
- Glauber-Gribov
  - Field theory formulation of Glauber model
  - Multiple collisions $\leftrightarrow$ Feynman diagrams
  - High energies: exchange of one or more Pomerons with one or more target nucleons (a closed string exchange)
- Formation zone (=materialization time)

Glauber Cascade

Quantum mechanical method to compute all relevant hadron-nucleus cross sections from hadron-nucleon scattering: $S_{hN}(\vec{b},s) = e^{i\chi_{hN}(\vec{b},s)} = \eta_{hN}(\vec{b},s)e^{2i\delta_{hN}(\vec{b},s)}$

and nuclear ground state wave function Ψ_i

Total $$\sigma_{hAT}(s) = 2\int d^2\vec{b}\int d^3\vec{u}\left|\Psi_i(\vec{u})\right|^2\left[1-\prod_{j=1}^{A}\mathrm{Re}\,S_{hN}(\vec{b}-\vec{r}_{j\perp},s)\right]$$

Elastic $$\sigma_{hAel}(s) = \int d^2\vec{b}\left|\int d^3\vec{u}\left|\Psi_i(\vec{u})\right|^2\left[1-\prod_{j=1}^{A}S_{hN}(\vec{b}-\vec{r}_{j\perp},s)\right]\right|^2$$

Scattering $$\sigma_{hA\Sigma f}(s) \equiv \sum_f \sigma_{hA\,fi}(s) = \int d^2\vec{b}\int d^3\vec{u}\left|\Psi_i(\vec{u})\right|^2\left\|\left[1-\prod_{j=1}^{A}S_{hN}(\vec{b}-\vec{r}_{j\perp},s)\right]\right\|^2$$

Absorption (particle prod.)

Absorption probability over a given *b* and nucleon configuration

$$\sigma_{hA\,abs}(s) \equiv \sigma_{hAT}(s) - \sigma_{hA\Sigma f}(s)$$

$$= \int d^2\vec{b}\int d^3\vec{u}\left|\Psi_i(\vec{u})\right|^2\left\{1-\left\{\prod_{j=1}^{A}1-\left[1-\left|S_{hN}(\vec{b}-\vec{r}_{j\perp},s)\right|^2\right]\right\}\right\}$$

Gribov interpretation of Glauber multiple collisions

Therefore the absorption cross section is just the integral in the impact parameter plane of the probability of getting at least one non-elastic hadron-nucleon collision

and the overall average number of collision is given by

$$\langle \nu \rangle = \frac{Z\sigma_{hpr} + N\sigma_{hnr}}{\sigma_{hAabs}}$$

- Glauber-Gribov model = Field theory formulation of Glauber model
- Multiple collision terms ⇒Feynman graphs
- At high energies : exchange of one or more pomerons with one or more target nucleons

- In the Dual Parton Model language: (neglecting higher order diagrams):

Interaction with ***n*** target nucleons ⇒ ***2n*** chains

- Two chains from projectile valence quarks + valence quarks of one target nucleon ⇒valence-valence chains
- 2(n-1) chains from sea quarks of the projectile + valence quarks of target nucleons ⇒2(n-1) sea-valence chains

Formation zone

Naively: "materialization" time (originally proposed by Stodolski).
Qualitative estimate:

In the frame where *p// =0*

$$\bar{t} = \Delta t \approx \frac{\hbar}{E_T} = \frac{\hbar}{\sqrt{p_T^2 + M^2}}$$

Particle proper time

$$\tau = \frac{M}{E_T}\bar{t} = \frac{\hbar M}{p_T^2 + M^2}$$

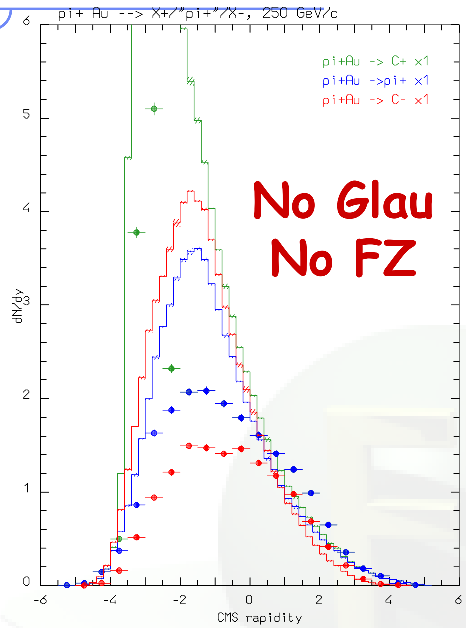
Going to the nucleus system

$$\Delta x_{for} \equiv \beta c \cdot t_{lab} \approx \frac{p_{lab}}{E_T}\bar{t} \approx \frac{p_{lab}}{M}\tau = k_{for}\frac{\hbar p_{lab}}{p_T^2 + M^2}$$

Condition for possible reinteraction inside a nucleus:

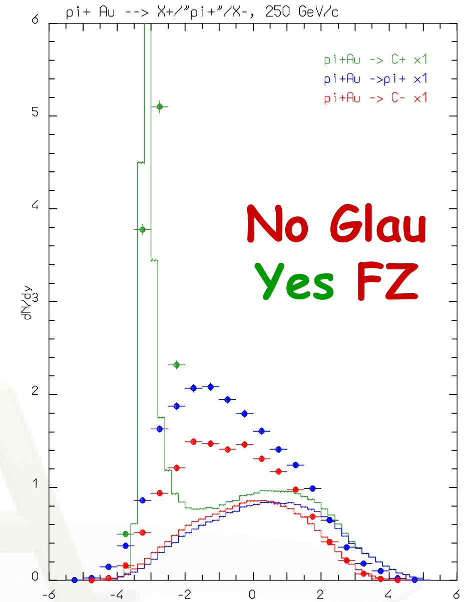
$$\Delta x_{for} \leq R_A \approx r_0 A^{\frac{1}{3}}$$

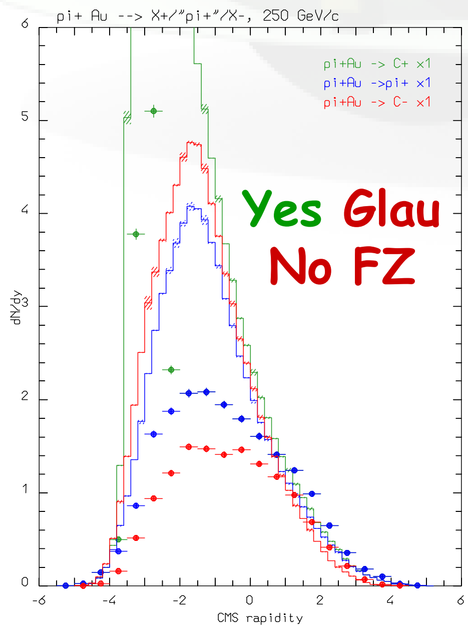
Effect of Glauber and Formation Zone

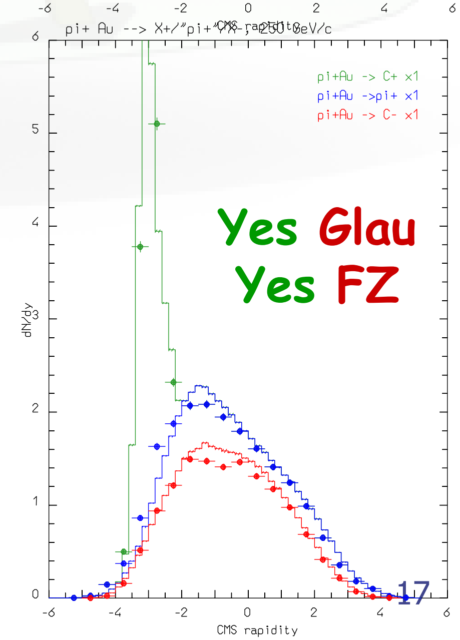


Rapidity distribution of charged particles produced in 250 GeV π^+ collisions on Gold

Points: exp. data (Agababyan et al., ZPC50, 361 (1991)).







WARNING

- PEANUT has been extended to cover the whole energy range in late 2006
- A less refined GINC model is available for projectile momenta above about 5 GeV/c, and was the only choice until recently
- However: the extended peanut is NOT yet the default, mainly because of some ambiguity in the definition of quasi-elastic cross sections.
- To activate PEANUT at all energies:

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PHYSICS 1000.  1000.   1000.   1000.   1000.   1000.   PEATHRESH
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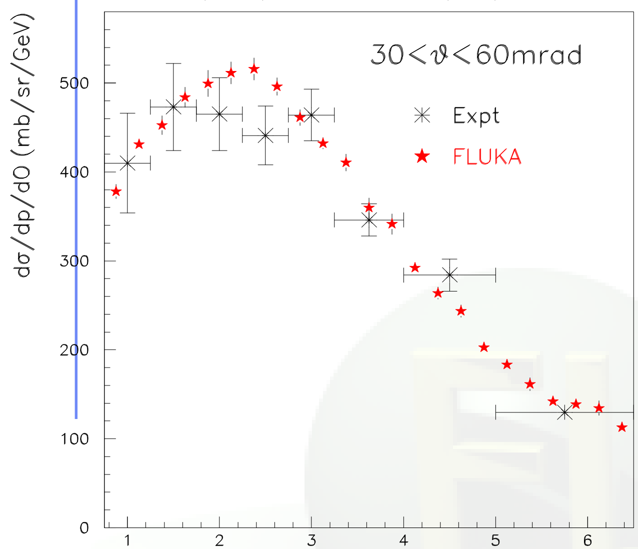
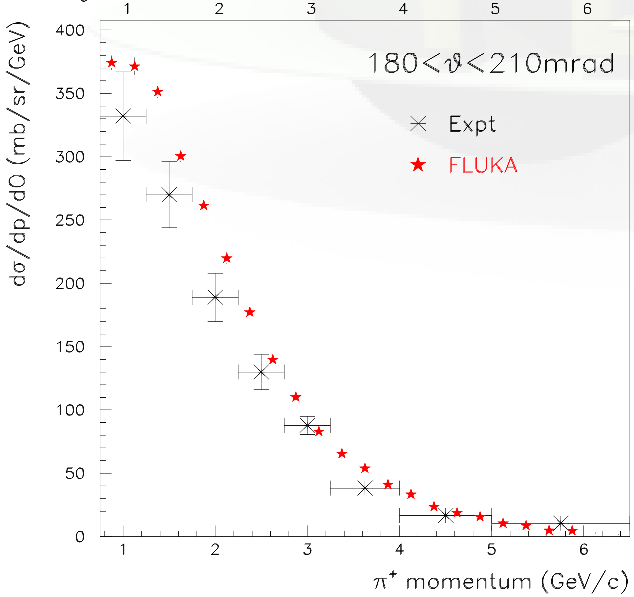
Nonelastic hA interactions at high energies: examples

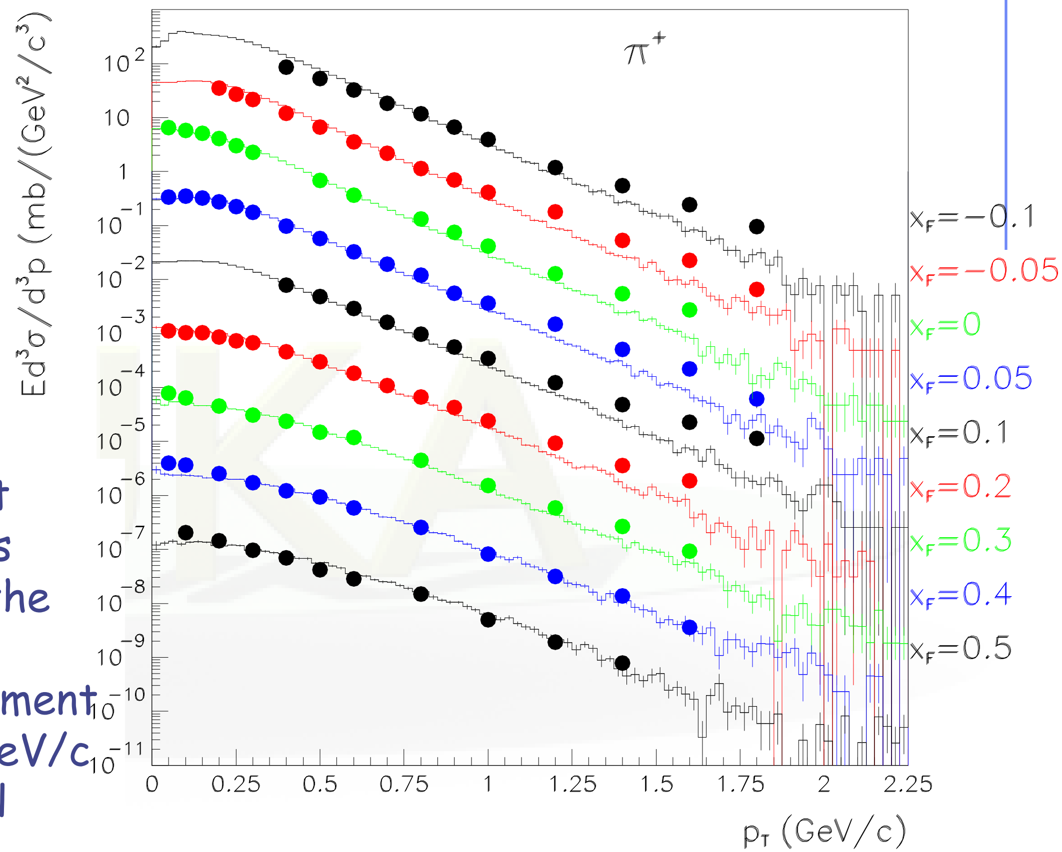
Harp expt. , 12.9 GeV/c p on Al





Recent results from the HARP experiment 12.9 GeV/c p on Al π^+ production at different angles

NA49 expt. , 158 GeV/c p on C



Double differential π+ production for p C interactions at 158 GeV/c, as measured by NA49 (symbols) and predicted by FLUKA (histograms)

Pions in nuclear medium

Pion-nucleon interactions: non-resonant + p-wave resonant Δ*'s.*

Δ in nuclear medium

decay

elastic scattering or charge exchange

reinteraction

pion absorption

$\rightarrow$ Δ width different from the free one

Assuming a Breit-Wigner for the free resonant cross section with width Γ_F

$$\sigma_{res}^{Free} = \frac{8\pi}{p_{cm}^2} \frac{M_\Delta^2 \Gamma_F (p_{cm})^2}{\left(s - M_\Delta^2\right)^2 + M_\Delta^2 \Gamma_F (p_{cm})^2}$$

Add "in medium" width (Oset et.al ,NPA 468, 631)

$$\frac{1}{2}\Gamma_T = \frac{1}{2}\Gamma_F - \mathrm{Im}\Sigma_\Delta, \quad \Sigma_\Delta = \Sigma_{qe} + \Sigma_2 + \Sigma_3$$

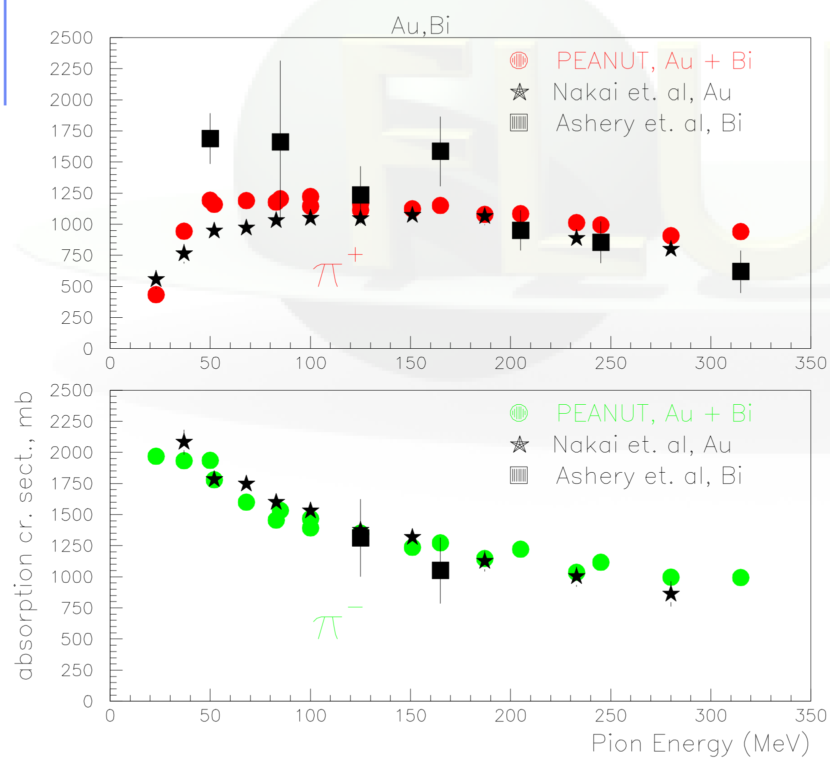
(Σ_{qe}*,* Σ_2*,* Σ_3 *= widths for quasielastic scattering, two and three body absorption)*

Add two-body s-wave absorption cross section from optical model

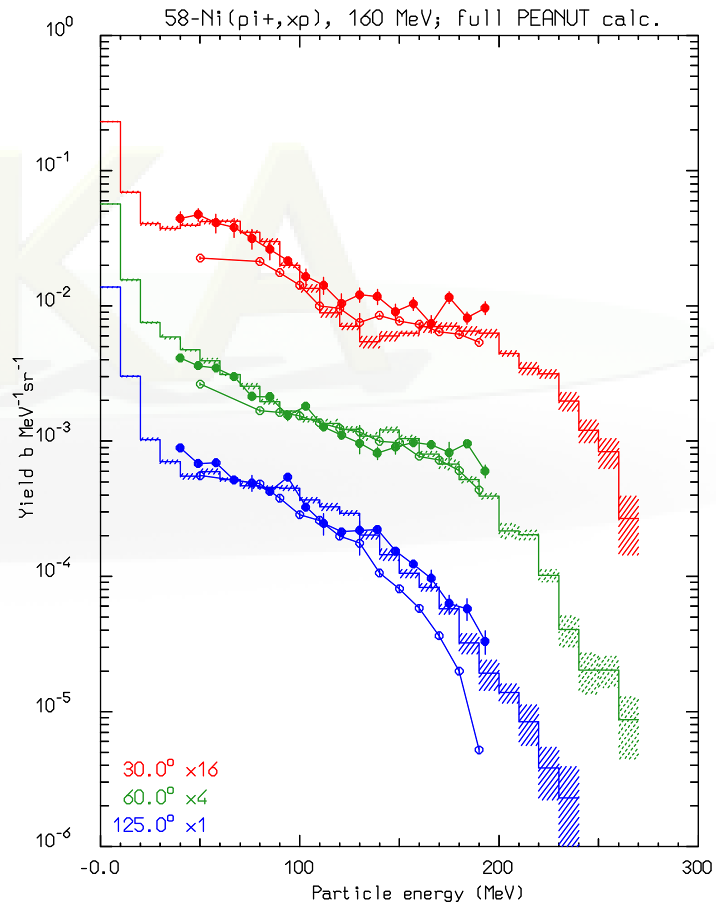
Nuclear potential for π: Energy dependent, resonant shape (+ Coulomb)

Pion absorption

Pion absorption cross section on Gold and Bismuth in the Δ resonance region (multibody absorption in PEANUT)



Emitted proton spectra at different angles, 160 MeV π^+ on ^{58}Ni
Phys. Rev. C41,2215 (1990)
Phys. Rev. C24,211 (1981)
Proton spectra extend up to 300 MeV



Preequilibrium emission

For E > π production threshold $\rightarrow$ only (G)INC models
At lower energies a variety of preequilibrium models

Two leading approaches

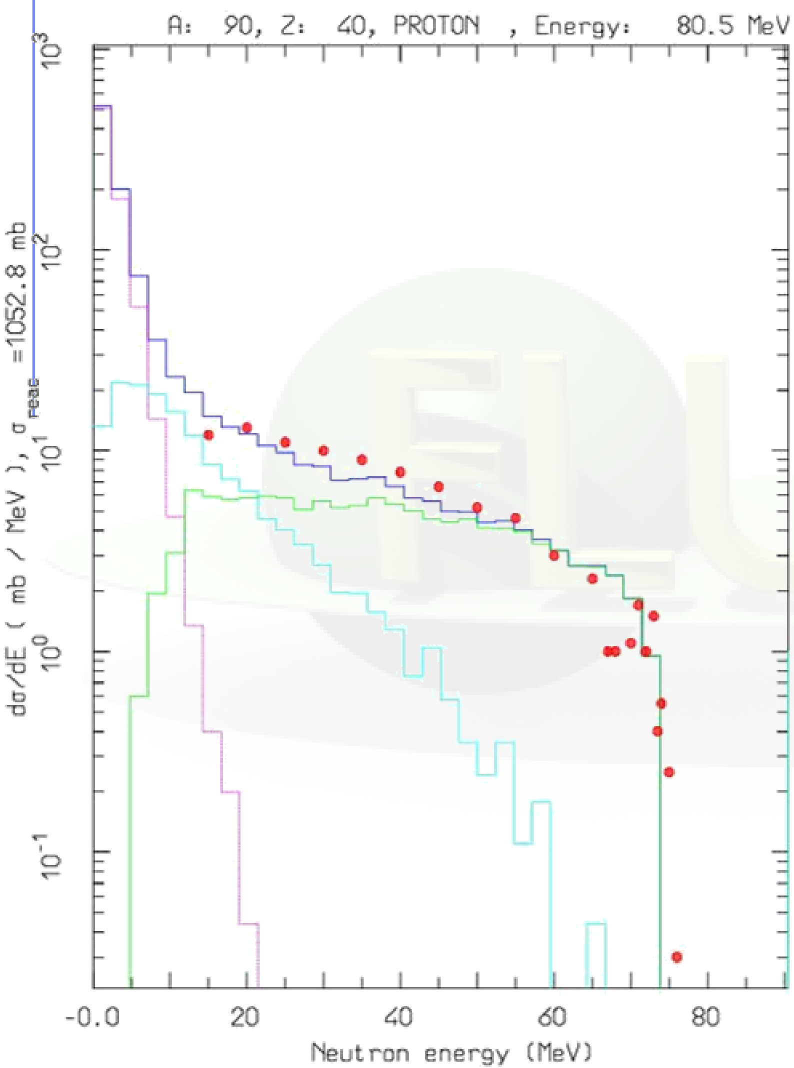
The quantum-mechanical multistep model:
Very good theoretical background
Complex, difficulties for multiple emissions

The semiclassical exciton model
Statistical assumptions
Simple and fast
Suitable for MC

Statistical assumption:
any partition of the excitation energy E^* among N, $N = N_h + N_p$, excitons has the same probability to occur
Step: nucleon-nucleon collision with $N_{n+1}=N_n+2$ ("never come back approximation)
Chain end = equilibrium = N_n sufficiently high or excitation energy below threshold

N_1 depends on the reaction type and cascade history

Thin target example



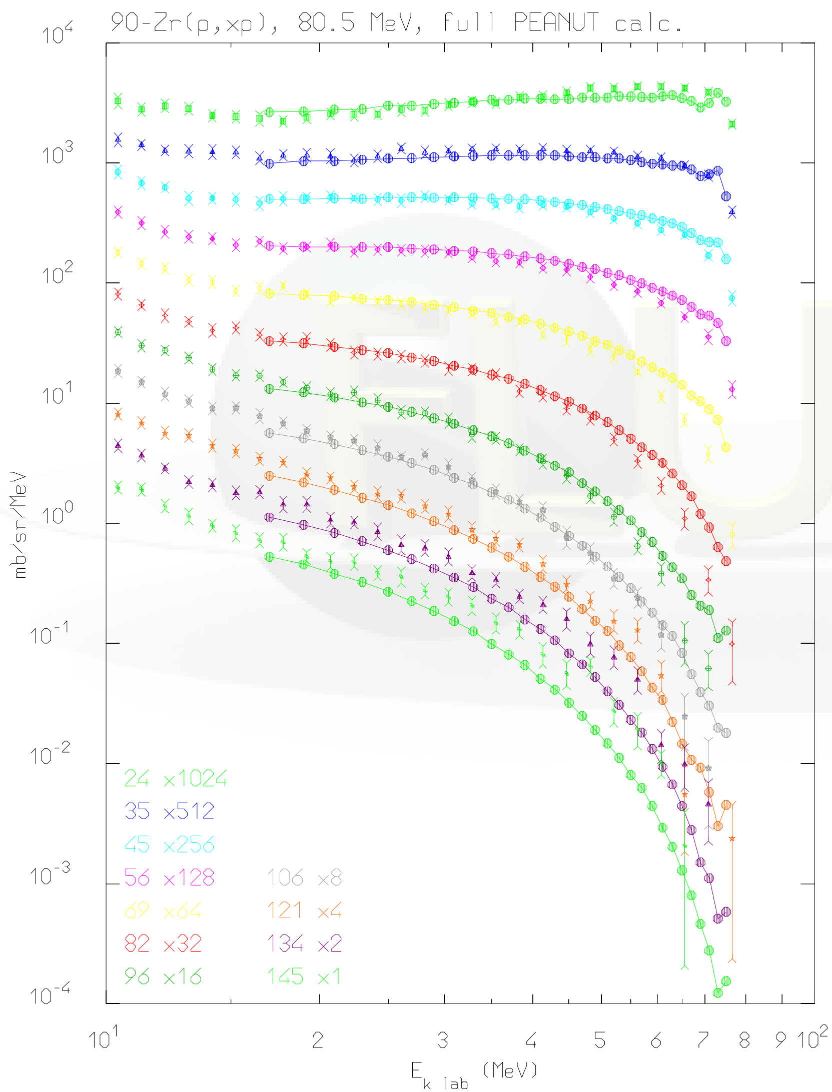
Angle-integrated $^{90}Zr(p,xn)$ at 80.5 MeV

The various lines show the total, INC, preequilibrium and evaporation contributions

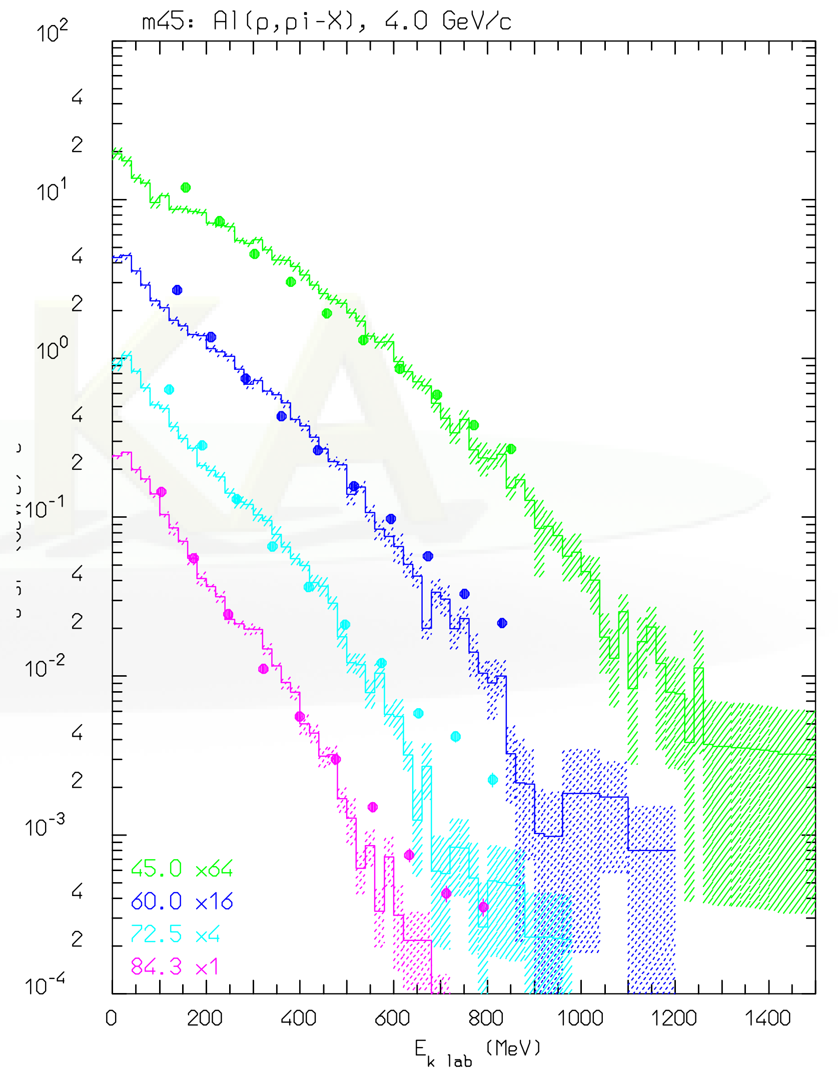
Experimental data from M. Trabandt et al., Phys. Rev. C39, 452 (1989)

Thin target examples

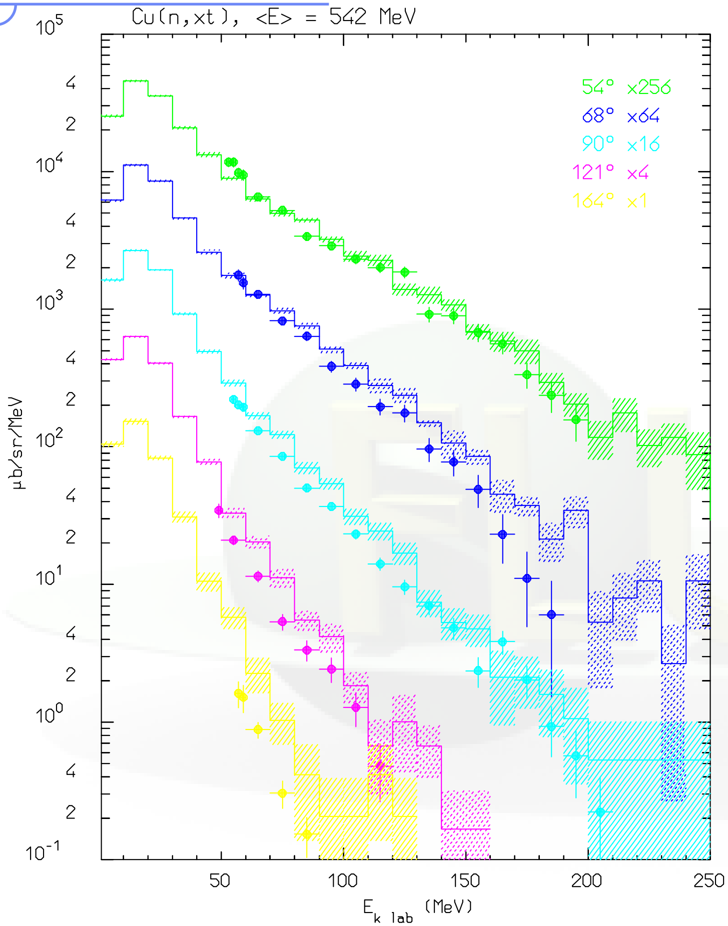
p + ^{90}Zr $\rightarrow$ p + X (80 MeV)



p + Al $\rightarrow$ π^- + X (4 GeV/c)



Coalescence



High energy light fragments are emitted through the coalescence mechanism: "put together" emitted nucleons that are near in phase space.

Example : double differential t production from 542 MeV neutrons on Copper

Warning: coalescence is OFF by default
Can be important, ex for . residual nuclei.
To activate it:

```
PHYSICS 1.                                              COALESCE
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If coalescence is on, switch on Heavy ion transport and interactions (see later)

Equilibrium particle emission (evaporation, fission and nuclear break-up)

From statistical considerations and the detailed balance principle, the probabilities for emitting a particle of mass m_j, spin $S_j\hbar$ and energy E, or of fissioning are given by:
(i, f for initial/final state, Fiss for fission saddle point)

Probability per unit time of emitting a particle j with energy E

$$P_J = \frac{(2S_j+1)m_j c}{\pi^2\hbar^3}\int_{V_j}^{U_i-Q_j-\Delta_f}\frac{\rho_f(U_f)}{\rho_i(U_i)}\sigma_{inv}(E)\,E\,dE$$

Probability per unit time of fissioning

$$P_{Fiss} = \frac{1}{2\pi\hbar}\int_0^{U_i-B_{Fiss}}\frac{\rho_{Fiss}(U_i-B_{Fiss}-E)}{\rho_i(U_i)}\,dE$$

- ρ's: nuclear level densities
- U's: excitation energies
- V_j's: possible Coulomb barrier for emitting a particle type j
- B_{Fiss}: fission barrier

- Q_j's: reaction Q for emitting a particle type j
- σ_{inv}: cross section for the inverse process
- Δ's: pairing energies

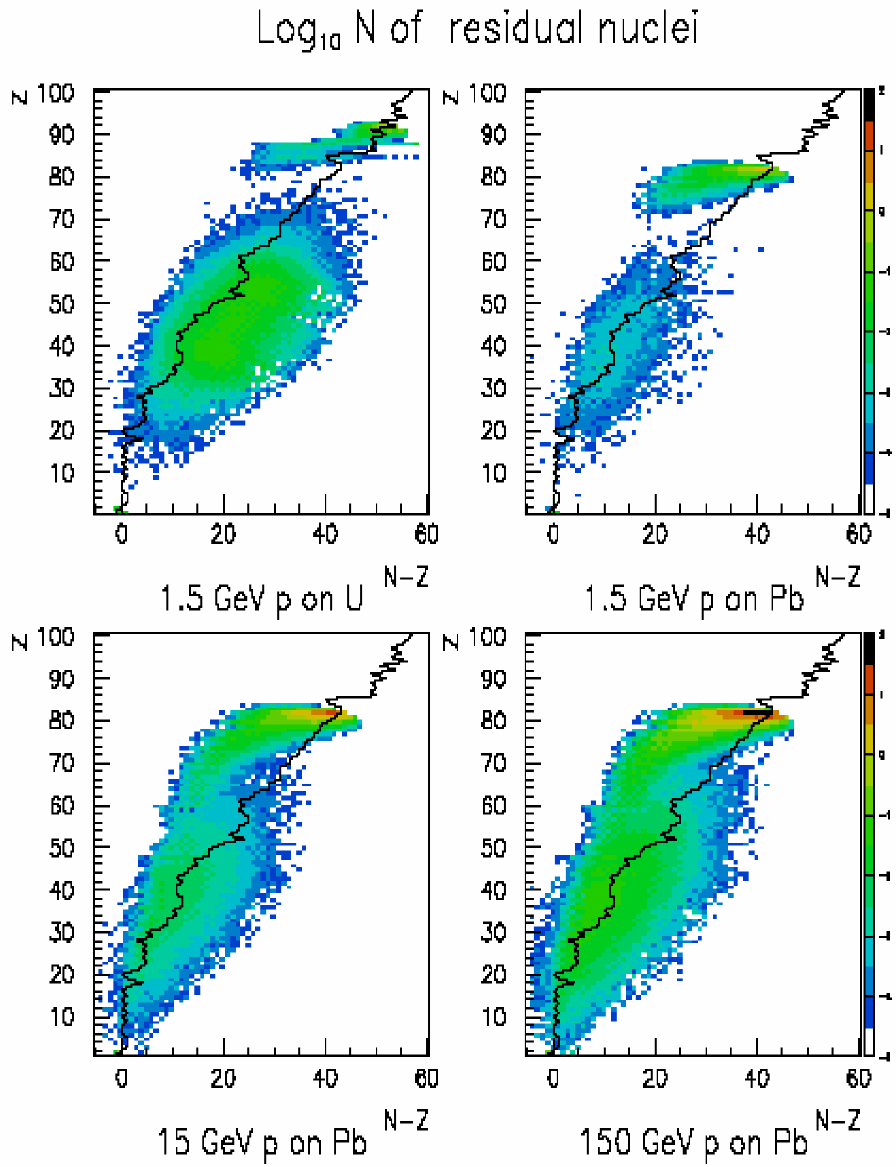
Neutron emission is strongly favoured because of the lack of any barrier
Heavy nuclei generally reach higher excitations because of more intense cascading

Equilibrium particle emission

- Evaporation: Weisskopf-Ewing approach
  - ~600 possible emitted particles/states (A<25) with an extended evaporation/fragmentation formalism
  - Full level density formula with level density parameter A,Z and excitation dependent
  - Inverse cross section with proper sub-barrier
  - Analytic solution for the emission widths (neglecting the level density dependence on U, taken into account by rejection
  - Emission energies from the width expression with no. approx.
- Fission: past, improved version of the Atchison algorithm, now
  - Γ_{fis} based of first principles, full competition with evaporation
  - Improved mass and charge widths
  - Myers and Swiatecki fission barriers, with exc. en. Dependent level density enhancement at saddle point
- Fermi Break-up for A<18 nuclei
  - ~ 50000 combinations included with up to 6 ejectiles
- γ de-excitation: statistical + rotational + tabulated levels

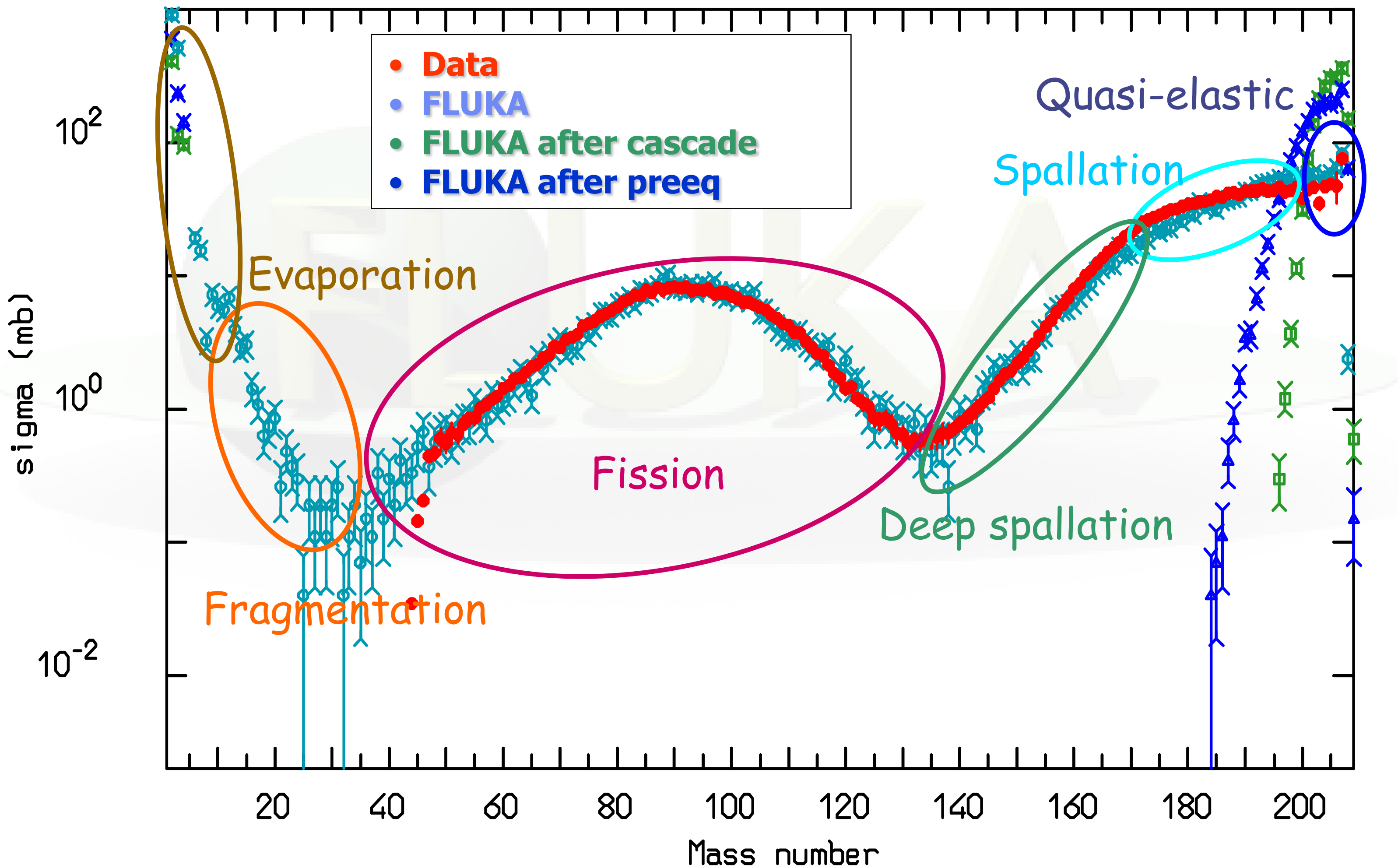
Residual Nuclei

- The production of residuals is the result of the last step of the nuclear reaction, thus it is influenced by all the previous stages
- Residual mass distributions are very well reproduced
- Residuals near to the compound mass are usually well reproduced
- However, the production of specific isotopes may be influenced by additional problems which have little or no impact on the emitted particle spectra (Sensitive to details of evaporation, Nuclear structure effects, Lack of spin-parity dependent calculations in most MC models)



Example of fission/evaporation

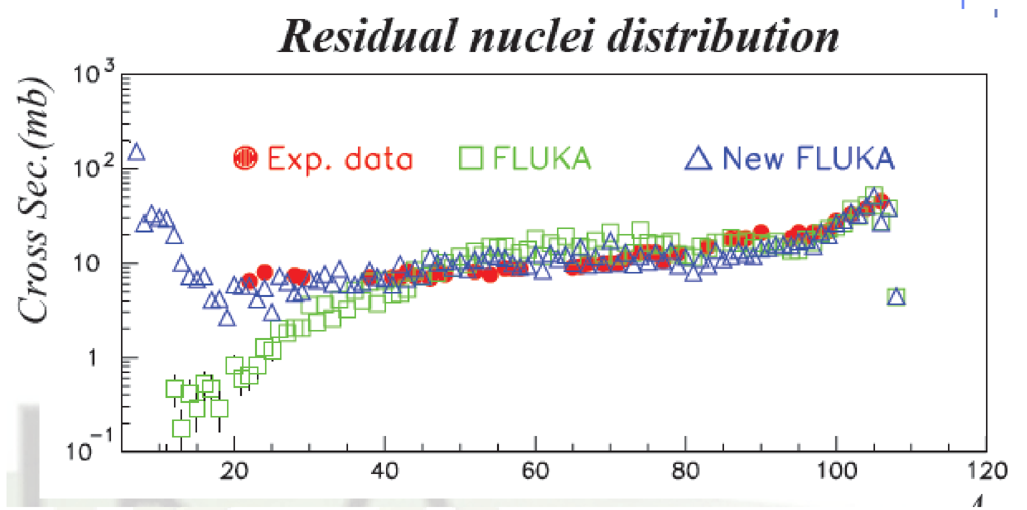
1 A GeV ^{208}Pb + p reactions Nucl. Phys. A 686 (2001) 481-524



Residual nuclei

Experimental and computed residual nuclei mass distribution for Ag(p,x)X at 300 GeV

Data from
Phys. Rev. C19 2388 (1979) and
Nucl. Phys. A543, 703 (1992)



The fragmentation model has much improved the FLUKA predictions

Also for A-A interactions

Warning: fragmentation is OFF by default, because it is a cpu-eater. It is NECESSARY to activate it for activation studies:

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PHYSICS 3.                                                  EVAPORAT
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If fragmentation is on, switch on Heavy ion transport and interactions (see later)